## London Theory Institute Lectures Series

## Bogdan Stefanski 6 D-branes and boundary states 9

## Informations

Pre-recorded Lectures : Youtube
Live Tutorial : Monday 7 th of December, 11h30


#### Abstract

D-branes are a remarkable set of soliton-like objects in string theory, which have played a critical role in many areas of string theory. Despite being non-perturbative objects in string theory, they can be understood exactly using world-sheet methods through the boundary state formalism. In this lecture, I introduce the notion of a boundary state and show how to use it to extract physical information about D-branes and the open strings that live on them.


## Exercices

Students are encouraged to try and solve the exercises by themselves and actively ask questions during the interactive live tutorial where solutions will be presented. The problems are :

1 Check that the momentum conjugate to $X$

$$
P \equiv \frac{\partial S}{\partial \dot{X}}
$$

satisfies the equal-time canonical commutation relation

$$
\left\{X\left(\sigma_{1}\right), P\left(\sigma_{2}\right)\right\}=i \delta\left(\sigma_{1}-\sigma_{2}\right)
$$

using the mode-expansions and commutation relations for $\alpha_{n}$ stated in the lecture.

2 Prove that for positive integers $n, k$

$$
\alpha_{n}\left(\alpha_{-n}\right)^{k}|0\rangle=n k\left(\alpha_{-n}\right)^{k-1}|0\rangle .
$$

3 Show that the Neumann boundary condition at $\tau=0$ implies that the Fourier modes of $X$ satisfy

$$
p_{0}=0, \quad \alpha_{-n}+\tilde{\alpha}_{n}=0
$$

and that the corresponding boundary state is

$$
|B\rangle=\mathrm{e}^{-\sum_{n=1}^{\infty} \frac{1}{n} \alpha_{-n} \tilde{\alpha}_{-n}}|B\rangle_{(0)}
$$

What is $|B\rangle_{(0)}$, the zero-mode part of the boundary state?
4 Expand the Born-Infeld action

$$
S_{B I}=T_{p} \int d^{p+1} \mathrm{xe}^{-\phi} \sqrt{-\operatorname{det}\left(g_{\mu \nu}+\mathcal{F}_{\mu \nu}\right)}
$$

to quadratic order in fields for a D3-brane, using the expansion

$$
g_{\mu \nu}=\eta_{\mu \nu}+\kappa h_{\mu \nu}
$$

where $\mu, \nu=0,1,2,3$. Take the antisymmetric $\mathcal{F}$ to have only non-zero components $\mathcal{F}_{01}=-\mathcal{F}_{10}=f_{1}$ and $\mathcal{F}_{23}=-\mathcal{F}_{23}=f_{2}$. Physically

$$
\mathcal{F} \equiv B+F
$$

where $B$ is the closed-string anti-symmetric NS-NS two form and $F$ is the field-strength for the $U(1)$ gauge field living on a single D -brane.
5 Prove that for positive integers $n, k$

$$
\alpha_{-n} \alpha_{n}\left(\alpha_{-n} \tilde{\alpha}_{-n}\right)^{k}|0\rangle=n k\left(\alpha_{-n} \tilde{\alpha}_{-n}\right)^{k}|0\rangle .
$$

6 All fields that enter the string world-sheet action transform under the Virasoro symmetry. Write down, in terms of Fourier modes, how the $(b, c)$ and $(\tilde{b}, \tilde{c})$ ghosts as well as the RNS fermions $\psi^{i}$ and $\tilde{\psi}^{i}$ contribute to $L_{n}$ and $\tilde{L}_{n}$. By using the symmetry preserving boundary conditions

$$
\left(L_{n}-\tilde{L}_{n}\right)|B\rangle=0
$$

deduce that the boundary conditions for the ghosts are

$$
c_{n}=-\tilde{c}_{-n}, \quad b_{n}=-\tilde{b}_{-n}
$$

and for fermions

$$
\psi_{r}^{\mu}=\mp i \tilde{\psi}_{-r}, \quad \psi_{r}^{a}= \pm \tilde{\psi}_{-r}^{a}
$$

for $\mu=0, \ldots, p$ and $a=p+1, \ldots, 9$, with $r$ half-integer/integer in the NS-NS/R-R sectors. Hint: To fix the relative sign between the $\mu$ and $a$ components you will need to consider the symmetric boundary conditions for super-Virasoro generators. What does the extra $\pm$ freedom correspond to for the fermions $\psi$ and $\tilde{\psi}$ ? Using these boundary conditions, write down the boundary state for ghosts and fermions. For this problem you may find this useful.

7 Show that the bosonic zero-mode contribution to the cylinder diagram for D-branes at positions $y_{1}$ and $y_{2}$ is given by

$$
{ }_{(0)}\left\langle B, y_{1}\right| e^{-2 \pi l H_{c}}\left|B, y_{2}\right\rangle_{(0)}=l^{(9-p) / 2} \mathrm{e}^{-\left(y_{1}-y_{2}\right)^{2} /(4 l \pi)} .
$$

Re-writing this in terms of the open string annulus modulus $t=\frac{1}{2 l}$ explain the meaning of this term in terms of open strings ending on the D-branes.
8 Calculate the cylinder diagram between a $\mathrm{D} p$-brane and a $\mathrm{D} q$-brane and show that in the superstring case this vanishes for $p+q=0 \bmod 4$.
9 Using a suitable orthogonal rotation, calculate the boundary state for a D-brane "at an angle" $\theta$ in the ( $X^{8}, X^{9}$ ) plane. Calculate the cylinder diagram for D -branes at angles to one another along two planes, i.e. with $\theta_{1}, \theta_{2}$. For what values of $\theta_{1}$ and $\theta_{2}$ does this vanish?

## References

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